

Electronic structure calculation methods using the Density Functional Theory-II

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Outline

- » Recap of HK theorems and KS equations
- » The physical meaning of the XC energy
- » Solution of a one-particle Schroedinger equation
- » Pseudo Potentials
- » The nitty-gritty of an actual KS calculation
 - » For molecules
 - » For solids
- » DFT software packages



The HK theorems

- » **HK Th 1:** *The external potential , and hence the total energy, is a unique functional of the electron density*
- » **HK Th 2:** *The density that minimizes the total energy is the exact ground state density.*



Kohn-Sham equations


If $\rho(\mathbf{r}) = \sum_{\lambda}^{\text{occupied}} |\psi_{\lambda}(\mathbf{r})|^2$ then

Minimizing $E[\rho] = T + E_{\text{ee}} + E_{\text{ext}}$ wrt ρ

leads to $H_{\text{KS}}[\rho] \psi_{\lambda} = \varepsilon_{\lambda} \psi_{\lambda}$ where

$$H_{\text{KS}}[\rho] = -\frac{\hbar^2 \nabla^2}{2m} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Coulomb}}[\rho(\mathbf{r})] + v_{\text{xc}}[\rho(\mathbf{r})]$$

Solve self – consistently


$$H_{\text{KS}}[\rho] \psi_{\lambda} = \varepsilon_{\lambda} \psi_{\lambda} ; \rho(\mathbf{r}) = \sum_{\lambda}^{\text{occupied}} |\psi_{\lambda}(\mathbf{r})|^2$$



What can be calculated from E_0 ?

- > Magnetization and magnetic susceptibility
- > EA, IP, Chemical potential, Band gap
- > Forces and Elastic constants
- > **Generalized Forces** and **Response Functions**,
which can be expressed as: $\frac{\partial E}{\partial A}$, $\frac{\partial^2 E}{\partial A \partial B}$
- > Using the KS wavefunctions, other properties can also be calculated...



Shortcomings of DFT

- » Overestimates binding energies (\approx by 15%)
- » Underestimates the gap (\approx by 40%)
- » “Good” for solids but “bad” for molecules

Overscreens, since functionals are taken from Jellium model

But **surprisingly successful** in predicting the correct crystal structure and elastic constants (<5% error) for a large variety of systems.



Improvements to DFT

- » **Generalized Gradients Approximation (GGA/MGGA)**: to handle non-uniformity in charge density (good for metals, still local):

$$E_{XC} = E_{XC}^{GGA} [\rho, |\nabla\rho|] \text{ or } E_{XC}^{MGGA} [\rho, |\nabla\rho|, \Delta\rho, \dots]$$

- » **Hybrid functionals** ($E_{xc} = a * E_{xc}[DFT] + b * E_x[HF]$) to remedy the bandgap problem
- » More advanced and general functionals

- > **non-local**

- > **orbital-dependent**



The physical meaning of the Exchange-correlation Energy

- » Pauli principle forbids two electrons to get too close to each other.
- » => Their Coulomb repulsion becomes weaker!
- » Coulomb interaction is overestimated by the electrostatic interaction ($E_{xc} < 0$)



Two-body terms

If $\rho(\mathbf{r}) = \langle \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \rangle$ and

$\rho(\mathbf{r}_1, \mathbf{r}_2) = \langle \sum_{i,j,i \neq j} \delta(\mathbf{r}_1 - \mathbf{r}_i) \delta(\mathbf{r}_2 - \mathbf{r}_j) \rangle$ then

$$\langle \sum_{ij,i \neq j} f(\mathbf{r}_i, \mathbf{r}_j) \rangle = \int d\mathbf{r}_1 d\mathbf{r}_2 \rho(\mathbf{r}_1, \mathbf{r}_2) f(\mathbf{r}_1, \mathbf{r}_2)$$

Self-interaction term



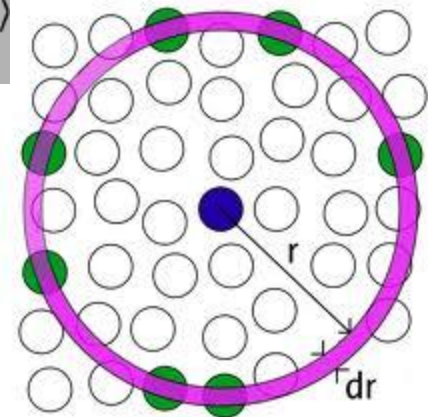
In the uncorrelated limit, $\rho(\mathbf{r}_1, \mathbf{r}_2) = \rho(\mathbf{r}_1)\rho(\mathbf{r}_2) - \rho(\mathbf{r}_2)\delta(\mathbf{r}_1 - \mathbf{r}_2)$

In the correlated limit,

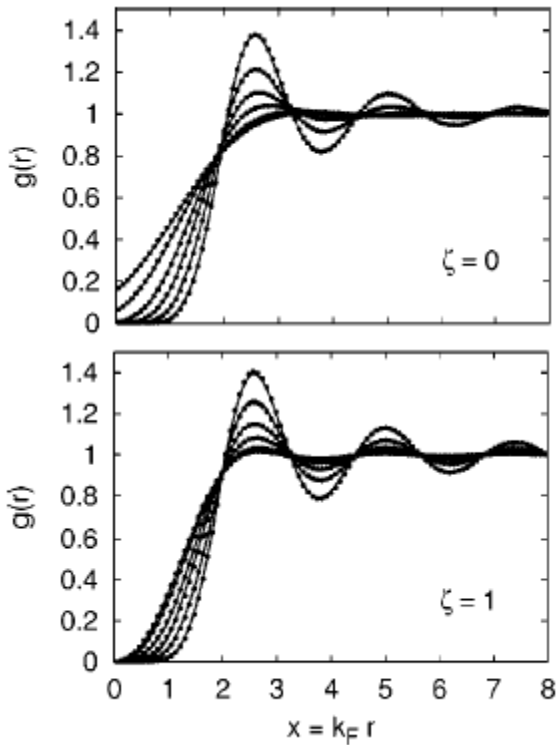
$$g(r) = \frac{V}{4\pi r^2 N^2} \langle \sum_i \sum_{j \neq i} \delta(r - r_{ij}) \rangle$$

$$\rho(\vec{\mathbf{r}}_1, \vec{\mathbf{r}}_2) = \rho(\vec{\mathbf{r}}_1) \rho(\vec{\mathbf{r}}_2) g(\vec{\mathbf{r}}_1, \vec{\mathbf{r}}_2)$$

g is called the pair correlation function and is only known for the Jellium model: $g_{\text{hom}}(\mathbf{r}, \rho)$



Electron correlations



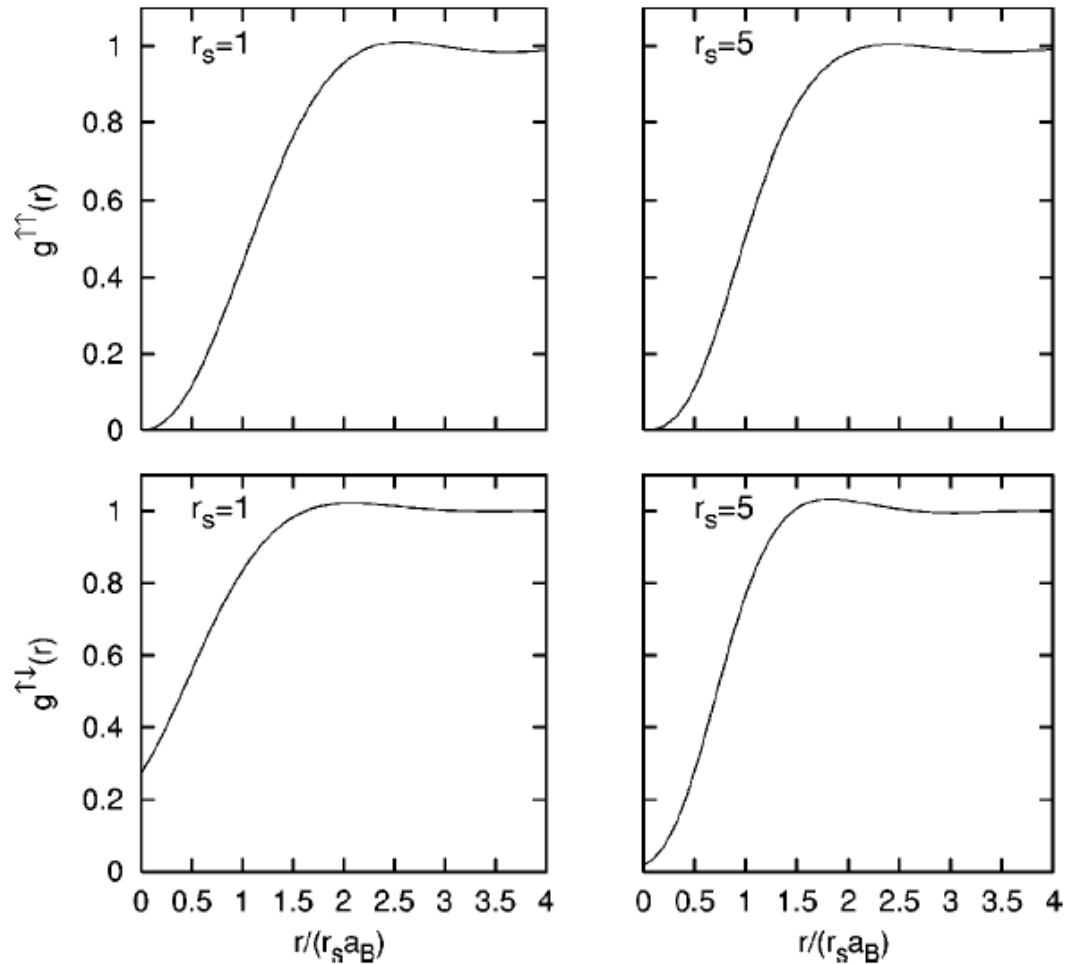
GORI-GIORGI, MORONI, AND BACHELET

PHYSICAL REVIEW B **70**, 115102 (2004)

2DEG For $r_s=1,2,5,10,20$

High density

Low density



B. DAVOUDI, M. POLINI, R. ASGARI, AND M. P. TOSI

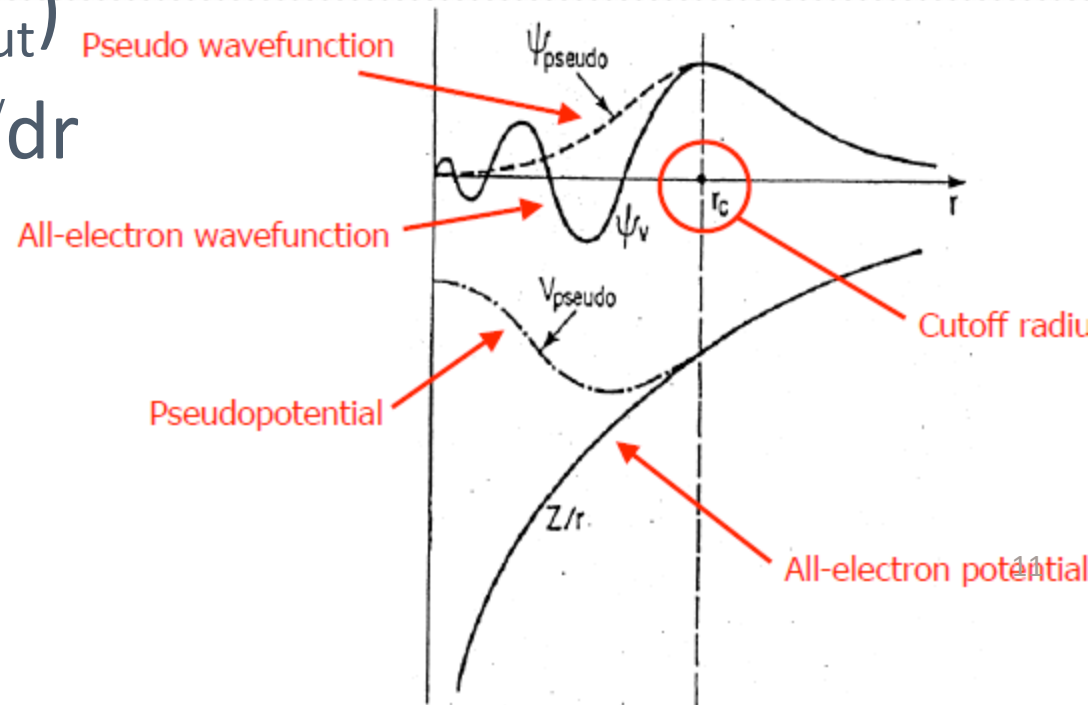
PHYSICAL REVIEW B **66**, 075110 (2002)

Pseudo-potentials

Core electrons do not contribute=>heavy calculations

Replace $(V_{ion} + \text{core electrons})$ by V_{pseudo}

- $\Psi_{pseudo} = \Psi_{exact} (r > r_{cut})$
- Scattering: $d \text{ Log } \psi_E / dr$
- Norm conservation
- Smoothness



How to solve a one-particle Schrodinger's equation?

» SPECTRAL METHODS:

» Consider the proper boundary conditions

» Choose accordingly a basis set ϕ_i :

$$H_{KS}[\rho] = \frac{p^2}{2m} + v_{\text{ext}}(r) + v_{\text{Coulomb}}[\rho(r)] + v_{\text{xc}}[\rho(r)]$$

$H|\psi_\lambda\rangle = \epsilon_\lambda|\psi_\lambda\rangle$; expand ψ_λ on a complete basis set:

$$|\psi_\lambda\rangle = \sum_j C_{\lambda j} |\phi_j\rangle$$

Multiplying the eigenvalue equation on the left by ϕ_i leads to:

$$\sum_j H_{ij} C_{\lambda j} = \epsilon_\lambda \sum_j S_{ij} C_{\lambda j}$$

with $H_{ij} = \langle \phi_i | H | \phi_j \rangle$ and $S_{ij} = \langle \phi_i | \phi_j \rangle$

How to solve a one-particle Schrodinger's equation?...

- » Construct the Hamiltonian and overlap **MATRIX**

$$\mathbf{H}_{ij} = \langle \phi_i | \mathbf{H}[\rho] | \phi_j \rangle; \mathbf{S}_{ij} = \langle \phi_i | \phi_j \rangle$$

- » Solve $\text{Det} [\mathbf{H} - \mathbf{E}\mathbf{S}] = 0$ (**Generalized eigenvalue problem**)
- » Get the density from the eigenstates:

$$\rho(\mathbf{r}) = \sum_{\lambda}^{\text{occupied}} |\psi_{\lambda}(\mathbf{r})|^2$$

- » Make it **self-consistent** since H depends itself on ρ

Flowchart of ESCalation algorithm

1) read coordinates of atoms from input file

2) Initialize the input charge density

3) Construct the Hamiltonian and overlap Matrix

4) Solve $D_{\text{eff}}[H+ES]=0$

Obtain the new charge density from the occupied eigenstates

Eventually mix the charge density with old ones

If converge calculate the ground state energy
else
go back to step 3 and use this density to calculate

Plane Wave basis for crystals

$$\varphi_k(r) = \frac{1}{\sqrt{\Omega}} e^{ikr}$$

$$\langle \varphi_k | H | \varphi_{k'} \rangle = \frac{\hbar^2 k^2}{2m} \delta_{k,k'} + V(G) \delta_{k-k',G}$$

k only couples to $k+nG \Rightarrow$ **G-periodicity** in k-space

\Rightarrow It is enough to keep k in the **First Brillouin Zone**

$$\langle \varphi_k | H | \varphi_{k+G} \rangle = V(G) ; \langle \varphi_k | H | \varphi_k \rangle = \frac{\hbar^2 k^2}{2m}$$

Number of plane waves decided by the choice of E_{cut} , usually leads to a large Hamiltonian matrix : $O(n_{\text{occ}} \times N_{\text{PW}}^2)$

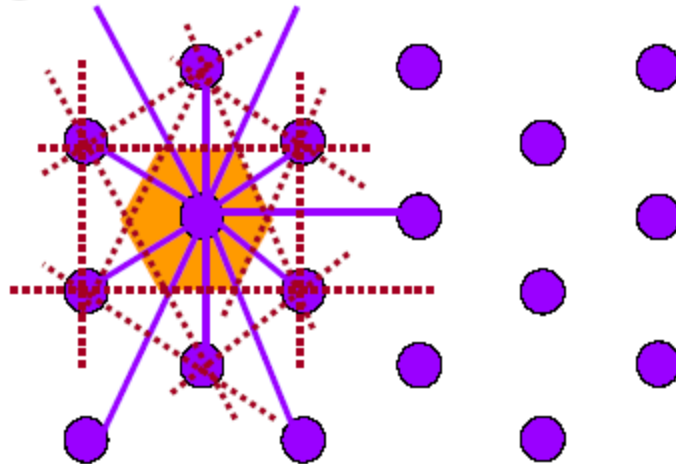
Structure of the Hamiltonian matrix in a PW scheme

$$\begin{array}{ccccccccc}
 G_1 & G_2 & G_3 & G_1 & G_2 & G_3 & G_1 & G_2 & G_3 \\
 \left[\begin{array}{ccc} H(k_1) & & \\ & 0 & \\ & & 0 \end{array} \right] & & & & & & & & \\
 & 0 & & \left[\begin{array}{ccc} H(k_2) & & \\ & 0 & \\ & & 0 \end{array} \right] & & & & & \\
 & & 0 & & & & \left[\begin{array}{ccc} H(k_3) & & \\ & 0 & \\ & & 0 \end{array} \right] & & \\
 \left[\begin{array}{ccc} & & \\ & & \\ & & \end{array} \right] & & & & & & & &
 \end{array}$$

Eigenvalues and eigenfunctions are periodic in the reciprocal space

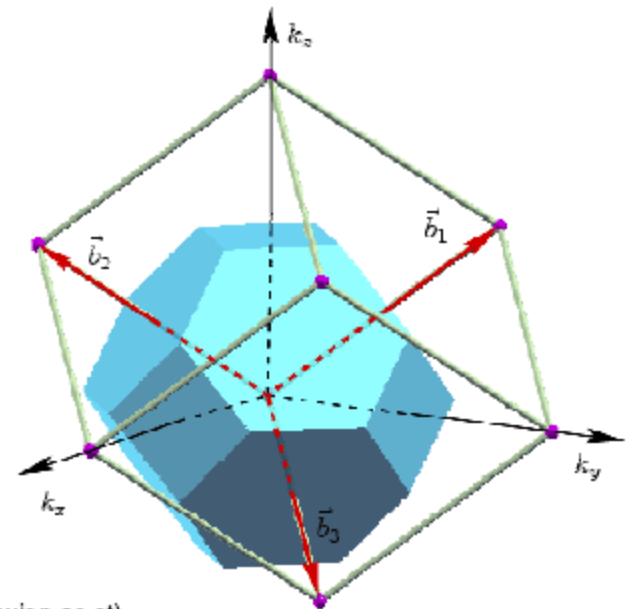
First Brillouin Zone

- Wigner-Seitz cell in reciprocal space.



(Or could choose to use parallelepiped defined by $\mathbf{b}_1, \mathbf{b}_2, \mathbf{b}_3$)

e.g., 1st BZ for FCC lattice →



(www.iue.tuwien.ac.at)



Atomic Orbital basis for crystals

$$\text{Bloch Basis : } \psi_{k\alpha}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}} e^{i\mathbf{k}\mathbf{R}} \phi_{\alpha}(\mathbf{r} - \mathbf{R})$$

Where α and β are atom or orbital indices within the unit cell $(\alpha, \beta) = \{1, \dots, n\}$

$$H_{\alpha\beta}(k) = \langle \psi_{k\alpha} | H | \psi_{k\beta} \rangle$$

is a rather small matrix: (n, n) for each k in the FBZ, where n is the number of orbitals per unit cell

Next Step: download, install and get familiar with a program

» Many Programs available:

- > Gaussian
- > SIESTA (free)
- > Wien2k
- > ELK (free)
- > VASP
- > ABINIT (free)
- > Quantum Espresso (free)
- > Check the Web

PW Calculation setup

» INPUTS:

- > Atomic coordinates
- > Unit cell
- > DFT functional (LDA, LSDA, GGA ...)
- > Cutoff Energy
- > Kpoints mesh
- > Diagonalization, mixing and Convergence threshold preferences

» 3 calculations required for a band structure & DOS

- > A SC with kpoints in the IBZ
- > A non-SC with chosen kpoints along symmetry directions (BS)
- > A finer kpoint mesh in the IBZ (DOS)

» OUTPUT:

- > Single total energy and forces on all atoms
- > Electronic eigenstates and eigenvalues, charge & spin density

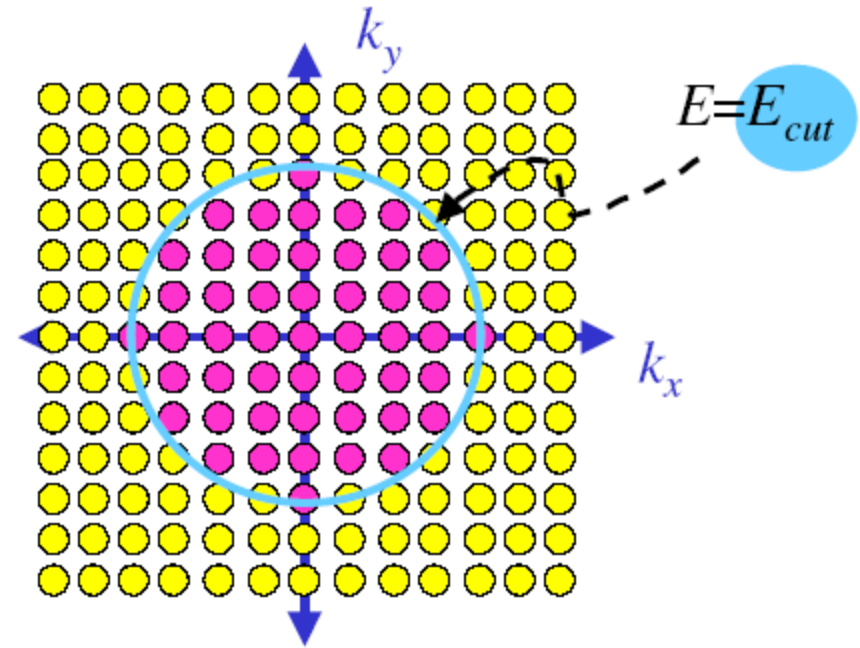


Truncating the Plane Wave Expansion

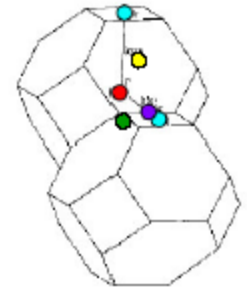
- In practice, the contribution from higher Fourier components (large $|\mathbf{k}+\mathbf{G}|$) is small.
- So **truncate** the expansion at some value of $|\mathbf{k}+\mathbf{G}|$.
- Traditional to express this **cut-off** in energy units:

$$\frac{\hbar^2 |\mathbf{k} + \mathbf{G}|^2}{2m} \leq E_{cut}$$

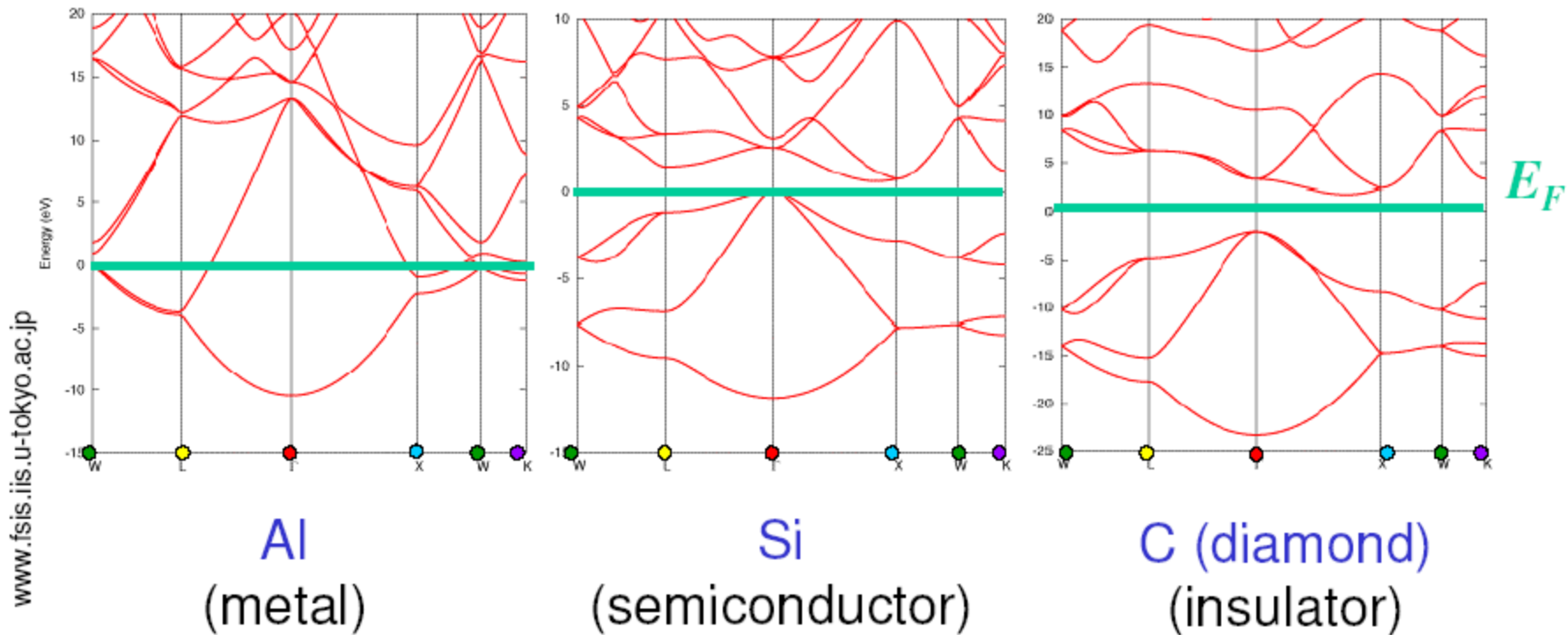
Input parameter `ecutwfc`



Band Structure

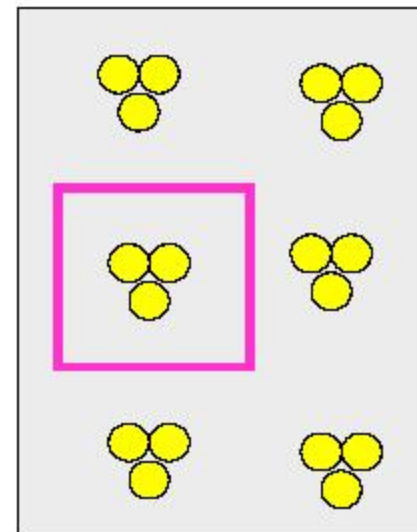
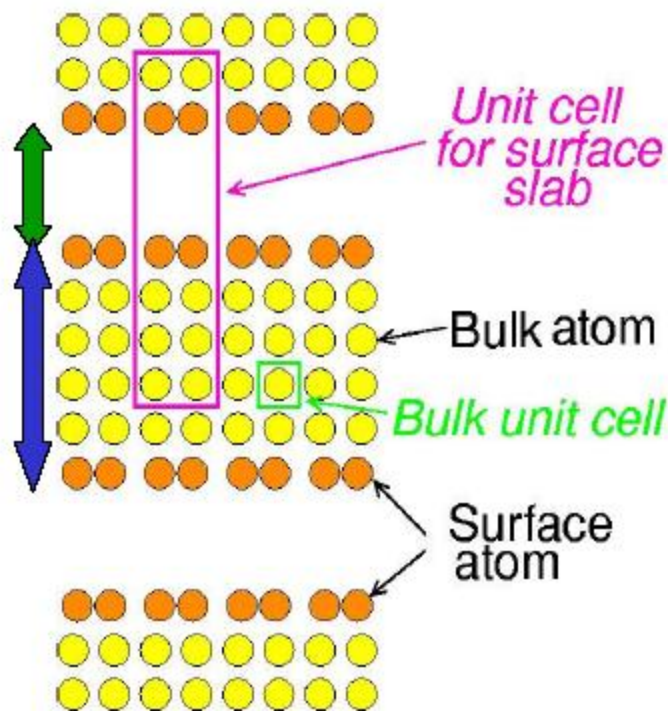
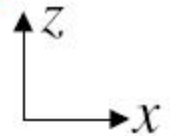


- “Band structure” (graph of ϵ_n vs. \mathbf{k}):
what states are available for electrons.
- At $T=0$, all states that lie below/above Fermi level are occupied/empty.



Artificially Periodic 'Supercells'

- Introduce 'vacuum' along the direction(s) where translational symmetry is broken to get **supercell**:



etc.



Brillouin Zone Sums

- Many quantities (e.g., density, total energy) involve integrals over \mathbf{k} :

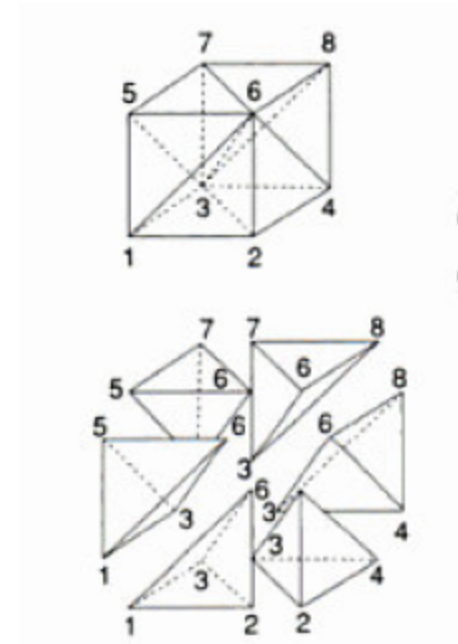
$$\langle P \rangle = \frac{\Omega}{(2\pi)^3} \sum_{n \text{ occ BZ}} \int P_n(\mathbf{k}) d^3k$$

- \mathbf{k} (wave-vector) is in the first Brillouin zone,
- n (band index) runs over occupied manifold.
- In principle, need infinite number of \mathbf{k} 's.
- In practice, sum over a finite number: BZ “Sampling”.



Interpolating between k-points

- In order to get sharp features in the DOS (e.g., van Hove singularities) correctly, use a method that interpolates between k-points.
- e.g., the **Linear Tetrahedron method**.
- Tetrahedra can be used to fill all space for any grid.
- Linear interpolation between values at vertices.



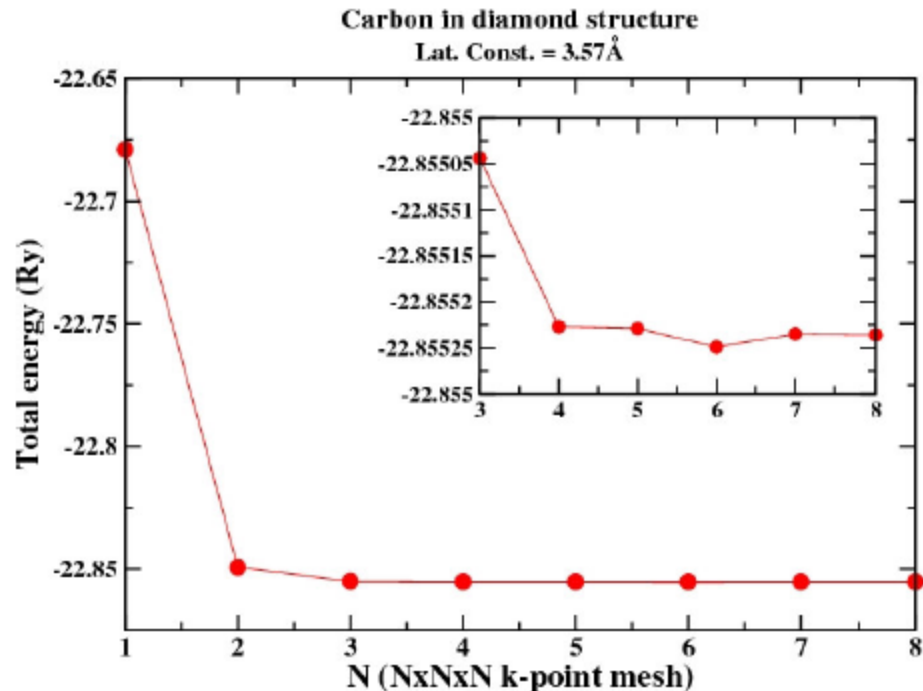
Bloch, Jepsen & Andersen



Shobhana Narasimhan, JNCASR

Keivan Estarjani, MIT nanoengineering group

Convergence wrt BZ sampling



Madhura Marathe

Note: Differences in energy usually converge faster than absolute value of total energy because of error cancellation (if supercells & k-points are identical or commensurate).



Shobhana Narasimhan, JNCASR

37

Convergence checks

- » Energy cutoff
- » Number of kpoints
- » Diagonalization
- » Self-consistency (mixing schemes)

Analyzing the output

- » Self-consistent electronic charge
- » Magnetization distribution
- » Total energy (stability)
- » Band structure or energy levels (approximate)

Using post-processing tools, can also compute :

- > Optical properties: dielectric matrix, optical conductivity
- > Mechanical properties: elastic constants
- > Chemical properties (hardness, IP, EA)
- > Phonons
- > Molecular dynamics: relaxation, reaction paths...